

MULTIFRACTAL ANALYSIS: EXTREME EVENTS CUT-OFF, CORRELATIONS, LINEARIZATION EFFECT.

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LINEARIZATION EFFECT

Multifractal Process $X(t)$ [?]:

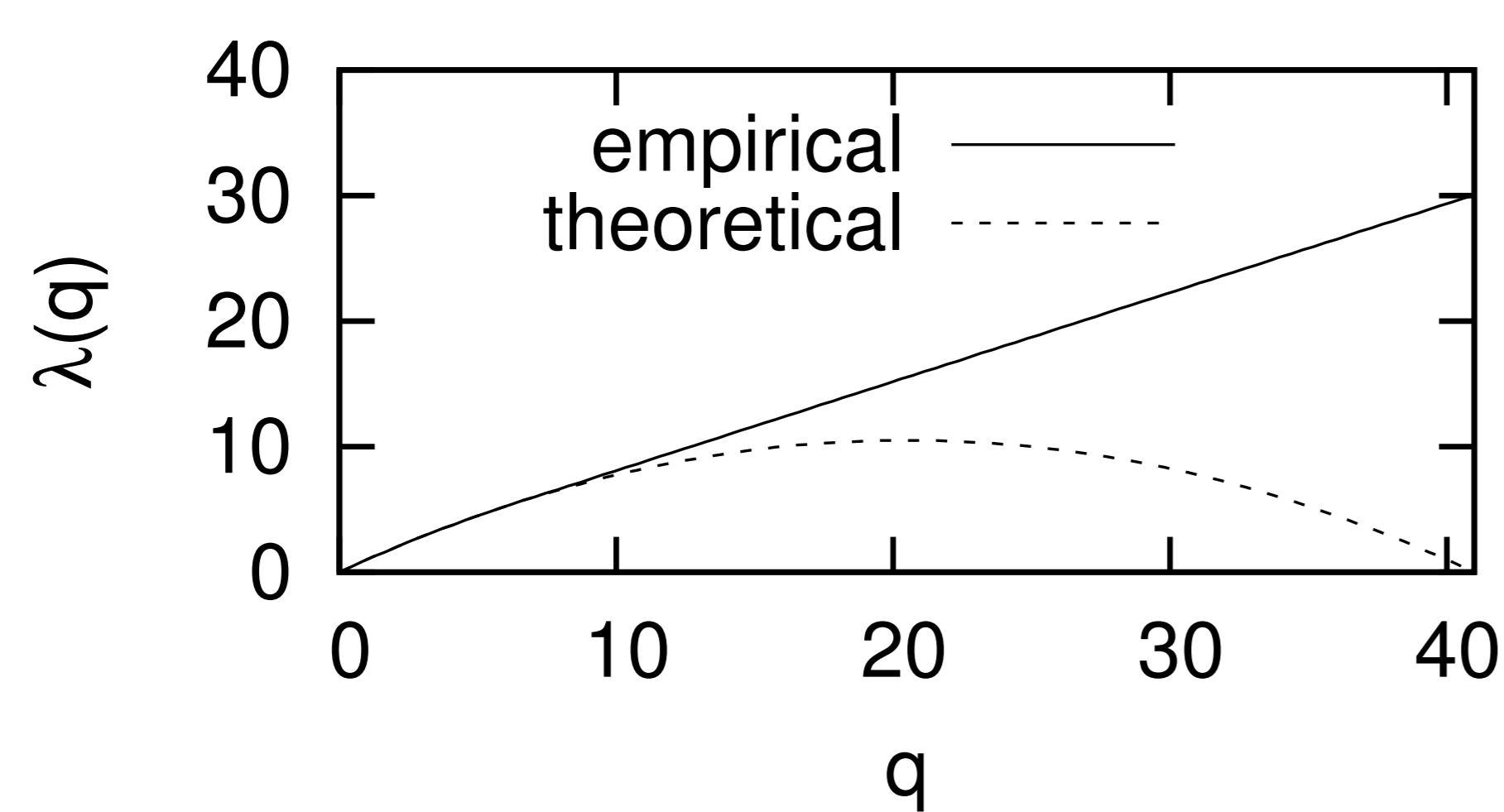
Increments $T(a, t) = X(t+a) - X(t)$,

$$\langle T_a^q \rangle = a^{\lambda(q)}, \quad a \rightarrow 0^+$$

$\lambda(q)$ is related to the Legendre transform of the multifractal spectrum

Moments estimator: $S_n(a, q) = \frac{1}{n} \sum_{k=1}^n T(a, t_k)^q$

Linearization effect [?] Beyond a critical moment q^* , experimental measurement using S_n diverges from $\lambda(q)$.



Large deviation properties

Large deviation theory [?] coupled with $\langle T_a^q \rangle = a^{\lambda(q)}$ indicates that the density probability p_a of $h_a = \frac{\ln T_a}{\ln a}$ verifies

$$p_a(h) \approx e^{-|\ln a| \psi(h)}$$

where $-\psi(h)$ is the Legendre transform of $\lambda(q)$.

EXTREME EVENTS AND CORRELATION [?]

Dominant moment contribution

Saddle point approximation,

$$\langle T(a, t)^q \rangle = \langle a^{qh_a(t)} \rangle \approx \int_{-\infty}^{+\infty} e^{-|\ln a| [qh + \psi(h)]} dh \approx e^{-|\ln a| [qh_m + \psi(h_m)]}$$

$$q + \psi'(h_m) = 0.$$

Long-range correlation : The number of effectively independent samples n_a scales as

$$n_a \approx \frac{L}{a}$$

Extreme events cut-off The finite size of the sample leads to a cut-off h^\dagger in the observable negative values of h_a :

$$P(h_a < h_a^\dagger) \approx \frac{1}{n_a}$$

Combining the expression of n_a and p_a leads to :

$$\psi(h^\dagger) = 1$$

Truncated moment A typical value of S_n can be evaluated by a truncated moment

$$S_n(a, q) \approx \int_{h^\dagger}^{+\infty} a^{qh} p_a(h) dh.$$

If $h_m(q) < h^\dagger$,

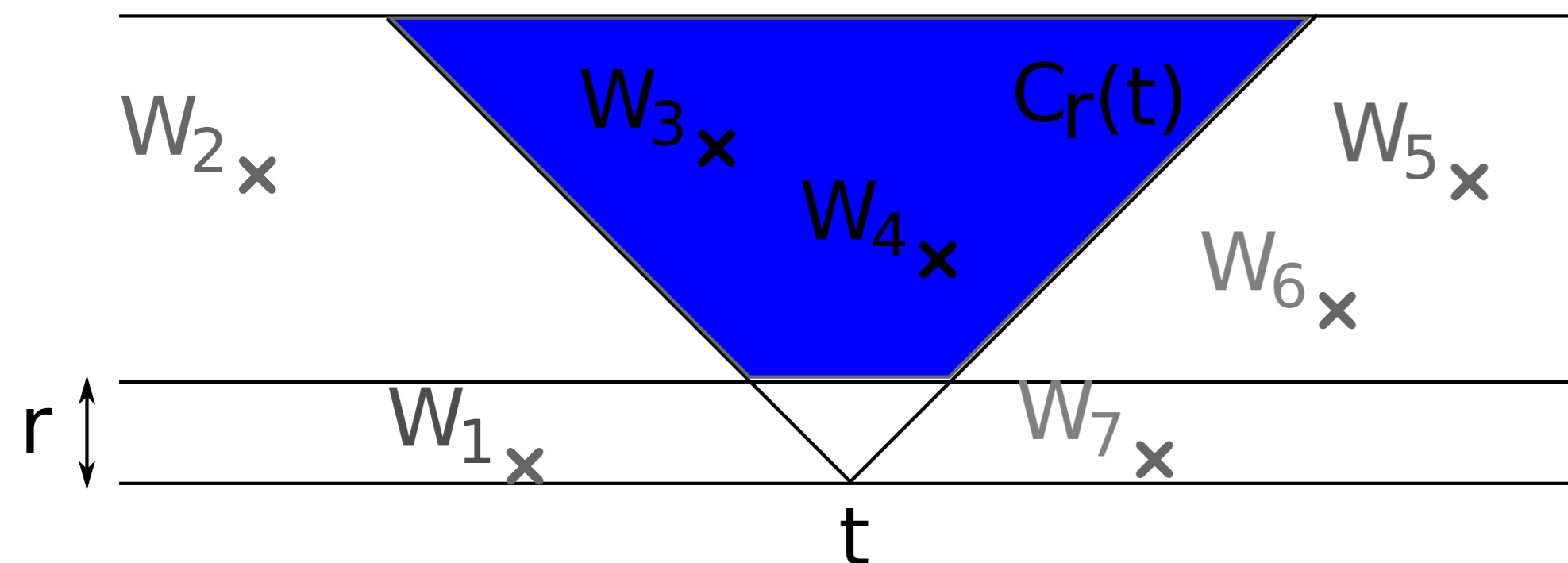
$$S_n(a, q) \approx a^{1+q\lambda(q^*)} \text{ (Linearization effect)}$$

Critical order : $h_m(q^*) = h^\dagger$

$$1 + q^* \lambda'(q^*) - \lambda(q^*) = 0.$$

COMPOUND POISSON PROCESSES [?]

Standard example of a multifractal process. Constructed from a Poisson point process W_n



Compound Poisson Cascade

$$Q_r(t) = B_r(t) \prod_{(t_i, r_i) \in \mathcal{C}_r(t)} W_i$$

where only multipliers associated with points belonging to the cone $\mathcal{C}_r(t)$ are taken into account and $B_r(t)$ is a normalising constant.

Compound Poisson Motion

$$X(t) = \lim_{r \rightarrow 0} \int_0^t Q_r(s) ds.$$

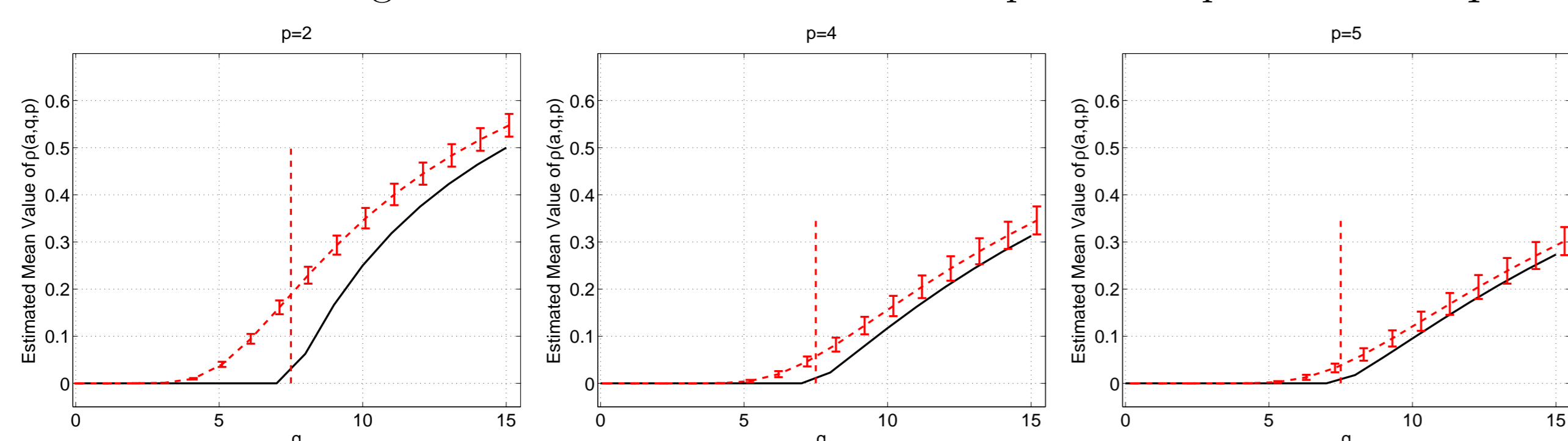
MONTE CARLO SIMULATION

Participation ratio

Classical tool for the study of glass transition.

$$\rho(a, q, p) = \frac{\sum_{k=1}^{n_a} |T(a, ka)|^{qp}}{(\sum_{k=1}^{n_a} |T(a, ka)|^q)^p}$$

Monte Carlo simulation gives a firm confirmation of the quantitative prediction of q^* .



Solid black line: $\langle \rho(a, q, p) \rangle$; dashed red line: $\langle \rho(a, q, p) \rangle_{MC}$ averaged over Monte-Carlo simulations, with 95% confidence intervals; red vertical dashed line: position of the critical q^*

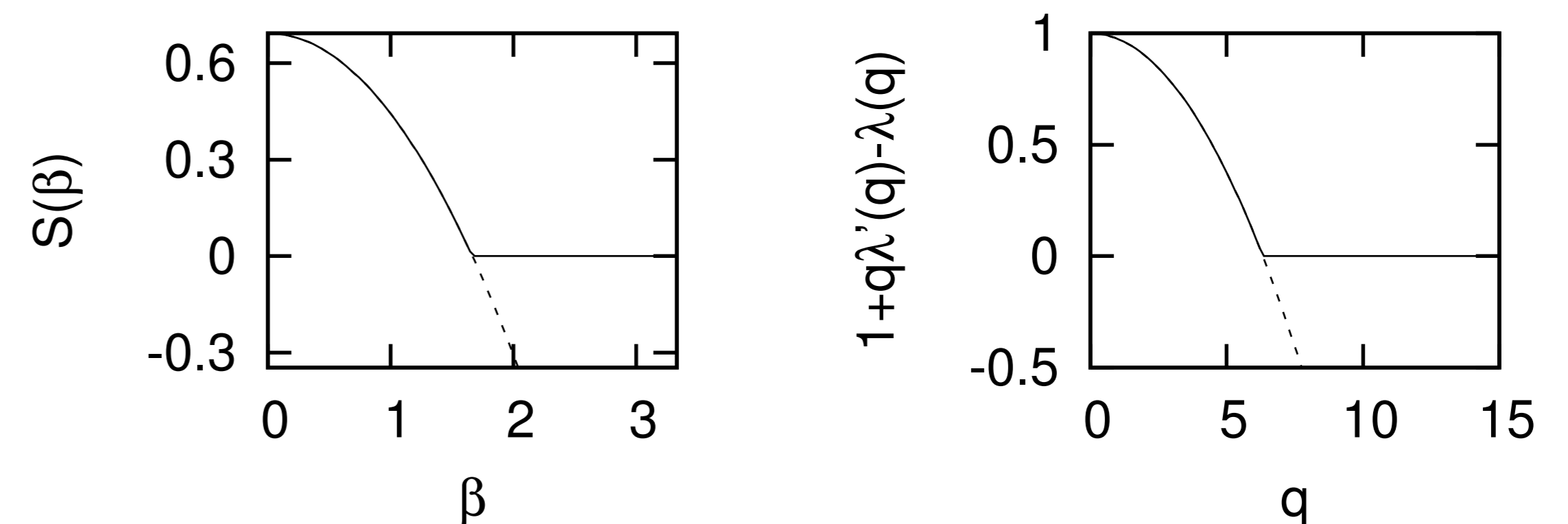
REM ANALOGY [?]

Formal analogy with the REM model

$$S_n(a, q) = \frac{1}{n} \sum_k T(a, t_k)^q = \frac{1}{n} \sum_k e^{-q |\ln a| h_a(t_k)} \approx Z$$

Z partition function of the REM. $\begin{cases} \text{inverse temperature } \beta \leftrightarrow q \\ \text{energy } E \leftrightarrow -|\ln a| h_a \end{cases}$

REM	Multifractal analysis
Entropy S	$\leftrightarrow 1 + q\lambda'(q) - \lambda(q)$
Glass transition at β_c	\leftrightarrow Linearization effect at q^* .



CONCLUSIONS

- A fruitful collaboration between statistical physics and multifractal analysis,
- Quantitative prediction of the Linearization effect threshold q^* ,
- Monte Carlo simulations confirm the similarity between the REM and the multifractal framework,
- q^* is an intrinsic characteristic of the process not a experimental limit,
- The Linearization effect is a glassy phase transition!

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